# Three-dimensional waves inside an open cavity and interactions with the impinging shear layer

Jérémy BASLEY, Julio SORIA, Luc R. PASTUR, and François LUSSEYRAN

## **1** Motivations

A separated flow over an open cavity (Fig.1) is primarily characterised by the enhancement of self-sustained oscillations [12]. *Kelvin-Helmholtz* travelling waves arise in the shear layer and *lock on* to the cavity length *L* due to an acoustic feedback loop [11, 13]. This leads to locked-on modes of oscillations often referred to as *Rossiter* frequencies in the compressible regime. When the external velocity is small with regards to sound speed ( $U_0 \ll c$ ), the corresponding frequencies  $f_n$  typically verify

$$f_n L/U_0 \approx n/2,\tag{1}$$

where  $n = \{1, 2, 3\}$  is the number of periods (or wavelengths) over *L*. Self-sustained oscillations represent highly energetic fluctuations, generating noise and drag and fluid-structure interactions. However, the envelope of those self-sustained oscillations is often disregarded, although drastic amplitude modulation is generally ob-

J. Soria

L. R. Pastur LIMSI/CNRS, BP 133, 91403 Orsay Cedex, FRANCE & Université Paris-Sud, F-91405 Orsay Cedex, FRANCE e-mail: luc.pastur@limsi.fr

F. Lusseyran

e-mail: francois.lusseyran@limsi.fr

J. Basley

Laboratoire d'Informatique pour la Mécanique et les Sciences de l'Ingénieur (LIMSI/CNRS), BP 133, 91403 Orsay Cedex, FRANCE & Université Paris-Sud, F-91405 Orsay Cedex, FRANCE e-mail: jeremy.basley@limsi.fr

Laboratory for Turbulence Research in Aerospace and Combustion (LTRAC), Department of Mechanical and Aerospace Engineering, Monash University, VIC 3800, AUSTRALIA e-mail: julio.soria@monash.edu.au

Laboratoire d'Informatique pour la Mécanique et les Sciences de l'Ingénieur (LIMSI/CNRS), BP 133, 91403 Orsay Cedex, FRANCE

served [6, 10]. An example of *very low* frequency modulations is shown in Fig.2. The shear layer locked-on mode constitutes the carrier while the envelope evolves over 20 times larger time-scales (about  $40 D/U_0$  time-steps).



**Fig. 2** Vorticity fluctuations  $\omega'_2 D/U_0$  out of Time-Resolved Particle Image Velocimetry data in a *xy*-plane, for L/D = 2.0 and  $L/\theta_0 = 82$  ( $U_0 = 1.38$  m/s). **Top** Time-series extracted near the impingement (x, y) = (1.5D, 0) and **bottom** picked-up snapshots.

This contribution aims at identifying the flow dynamics responsible for such low frequencies. The three-dimensional (3D) organisation of the flow is believed to be the source of the amplitude modulation [10]. In particular, the extensive literature on the recirculating flow inside the cavity has shown how *centrifugal instabilities* lead to slow 3D waves. In both lid-driven [1, 16] and shear-driven cavity flows [7, 8, 9], vortical structures are observed along the span of the cavity. They are typically associated with low frequencies  $f_{ci}$  such that

$$f_{ci}D/U_0 \leqslant 0.05,\tag{2}$$

consistent with the time scales of the envelope observed in Fig.2. Furthermore, Brès & Colonius [5] have already pointed out such low frequencies co-existing with *Rossiter* oscillations of the shear layer, in direct numerical simulations (DNS) of compressible open cavity flows.

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## 2 Experiments

The present study deals with an experimental cavity flow in the incompressible regime. Consider a laminar incoming flow, for which  $\theta_0$  is the momentum thickness of the Blasius boundary layer at separation. The dimensionless cavity length  $L/\theta_0$  and the Reynolds number  $\text{Re}_D = U_0 D/\nu$  are primary control parameters. Two campaigns have been conducted to encompass a wide range of the parameter space:  $1500 \leq \text{Re}_D \leq 9000$  and  $23 \leq L/\theta_0 \leq 104$ . The space-time evolution of the 3D flow dynamics is investigated through Time-Resolved Particle Image Velocimetry (TRPIV) measurements performed in two planes. On one hand, wind-tunnel experiments at LIMSI focus on the shear layer waves and the primary dynamics. Highspeed PIV measurements (500Hz) are performed in a xy-plane, streamwise normal to the bottom of the cavity [2, 3, 4]. On the other hand, the water-tunnel campaign at LTRAC is concerned with the spanwise extension of the flow [2]. It consists of high-resolution PIV data out of a zx-plane, parallel to the bottom of the cavity at y/D = -0.1. Velocity fields are obtained from particle images recorded by three synchronised cameras and processed using a cross-correlation algorithm [15]. Cavity span is larger at LTRAC (S = 10D) than at LIMSI (S = 6D) so as to better identify spanwise wavelengths. Note that the spanwise extension is large enough in both cases to ensure that influence of the end-walls remains secondary to intrinsic stability properties of the recirculating inner-flow.

## **3** Slow dynamics inside the cavity

Since PIV data are time-resolved, spectral filtering can be performed to separate the low frequency range from the dominant locked-on frequency and its harmonics. Such a decomposition is illustrated in Fig.3. The locked-on frequency of the shear layer  $f_a$  is such that  $St_a = f_a D/U_0 = 0.49$ , while the local maximum of the low frequency range is observed at  $St_\Delta = f_\Delta D/U_0 = 0.024$ .



**Fig. 3** Filtered vorticity fluctuations in a *xy*-plane (LIMSI), for  $L/D = 2.0 \& L/\theta_0 = 82$ . Examples of filtered snapshots, **left**, around shear layer frequencies, **right**, for low frequencies only.

As expected, the dynamics corresponding to  $f_a$  is restricted to the shear layer and to inflows along the downstream wall of the cavity. On the contrary, low frequency dynamics encompasses the entire inner-flow. More particularly, slow dynamics organises as large-scale fluctuations winding up around the main recirculation and implying outflows near the impingement. Such a spatial structure associated with frequencies matching Eq. 2 suggests 3D waves resulting from centrifugal instabilities, as observed in numerical studies by [5, 9]. The identification of those waves requires an investigation of the spanwise extension of the flow.

### **4** Spanwise travelling waves

The spanwise extension of the inner-flow is investigated through PIV measurements in a *zx*-plane at y/D = -0.1. Since recordings are time-resolved, Fourier transform can be applied to the time-series at each point of the velocity fields [14, 4]. Spatial structures associated with a single frequency are hence identified. They are referred to as *global Fourier modes* in the following.

The most energetic global Fourier modes are always associated with Strouhal numbers matching  $St_{\Delta} = f_{\Delta}D/U_0 \approx 0.02$ , as shown in Fig.4. The salient dynamics of the inner-flow organise as planar spanwise-travelling waves. Highly coherent spanwise oscillations appear for large areas of the *zx*-plane, yielding a unique spanwise wavelength  $\lambda \approx D$ . In other words, the dominant 3D dynamics of the inner-flow can be represented by *monochromatic spanwise travelling waves*. Such space-time feature typically come by pair of counter-propagating waves, which may exclude each other or partially overlap. In the latter case, interference leads locally to a (quasi) *standing wave* [2].

With increasing control parameters, shear layer disturbances get stronger and make the inner-flow more unsteady. However, the most salient modes connected to slow dynamics persist as monochromatic spanwise travelling waves. This is shown in [3], for  $L/\theta_0 = 76$ , Re<sub>D</sub> = 6800. Hence, it is reasonable to assume that 3D dynamics corresponding to frequency St<sub> $\Delta$ </sub> can write in the (dimensionless) form

$$\psi_{\Lambda}(x, y, z, t) = \zeta_{\Delta}(x, y) \times \exp[i(\beta z - 2\pi \operatorname{St}_{\Delta} t)], \qquad (3)$$

where  $\beta$  is the spanwise wavenumber and  $\zeta_{\Delta}(x, y)$  stands for the global Fourier mode associated with St<sub> $\Delta$ </sub> in a *xy*-plane. With that assumption, the spanwise dynamics of  $\psi_{\Delta}$  is equivalent to the temporal dynamics in a single *xy*-plane. Through a simple dispersion relationship, it comes

$$\frac{\partial \psi_{\Delta}}{\partial z} \propto \frac{\partial \psi_{\Delta}}{\partial t} \qquad \left(\frac{\partial \psi_{\Delta}}{\partial z} = \beta \,\psi_{\Delta} = \frac{2\pi \mathrm{St}_{\Delta}}{c_{\Delta}} \,\psi_{\Delta} = -\frac{1}{c_{\Delta}} \frac{\partial \psi_{\Delta}}{\partial t}\right), \tag{4}$$

with  $c_{\Delta}$  the (constant) spanwise phase velocity of the wave. Three-dimensional dynamics associated with frequency St<sub> $\Delta$ </sub> can therefore be estimated through a reconstruction in the space-time volume (*x*, *y*, *t*) using the *xy*-plane Fourier mode  $\zeta_{\Delta}(x, y)$ (Fig.5). Time-wise axis here stands for the spanwise extension of the flow.



#### **5** Conclusion and outlook

 $\frac{1}{x/D}$ 

We have shown that the lowest amplitude modulations of the shear layer waves are related to centrifugal instabilities inside the cavity. In fact, inner-flow fluctuations can be primarily modelled as monochromatic spanwise travelling waves of frequency  $f_{\Delta}D/U_0 \approx 0.02$ . Such a model implies spanwise derivatives become proportional to temporal derivatives. As a result, 3D structures associated with  $f_{\Delta}$  can be estimated out of the two-dimensional space-time dynamics at a given position *z*. Future works will aim to characterise more precisely the non-linear interactions between shear layer waves and 3D slow dynamics inside the cavity.

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